



# Problems in solid Earth deformation: crust and upper mantle

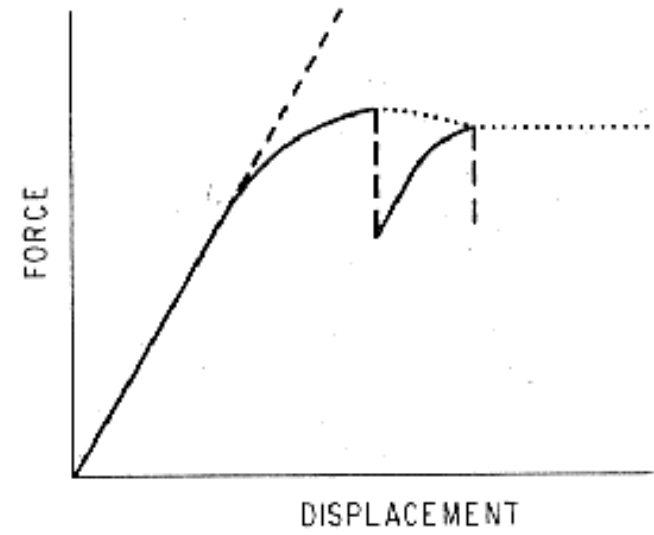
Yuri Fialko

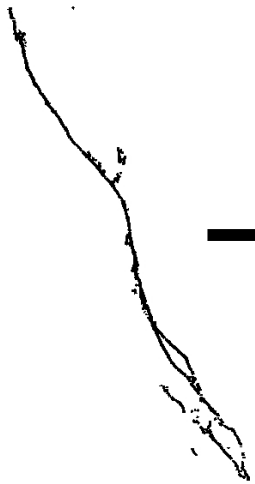
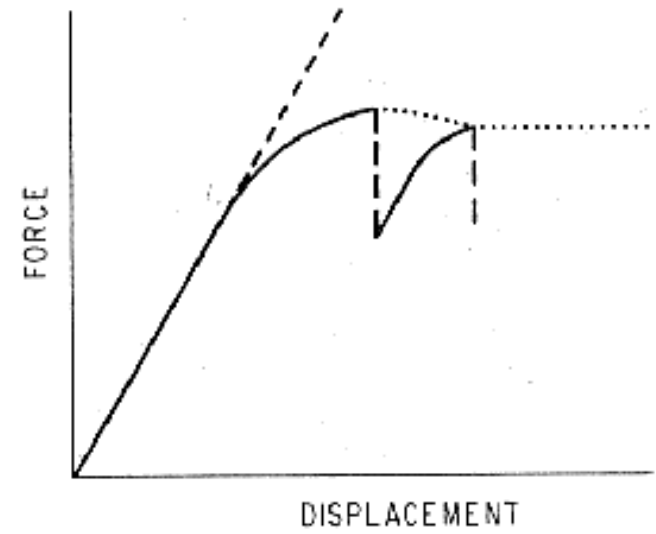
S. Barbot  
J. Pearce  
Y. Hamiel

Institute of Geophysics and Planetary Physics  
Scripps Institution of Oceanography  
University of California San Diego

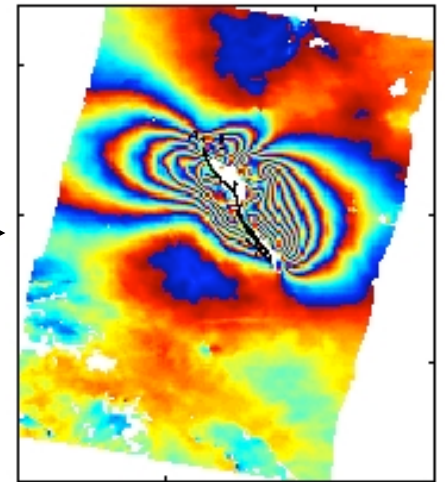
# Overview:

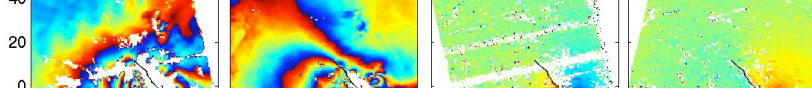
- Data-driven models
- Quest for “realistic” constitutive relationships for lithosphere, including upper (brittle) crust, lower crust, and upper mantle
- Data *require* material heterogeneity, non-linear rheologies, localization
  - models need to be sufficiently flexible to resolve multiple spatial and temporal scales
  - sufficient flexibility/efficiency in generating many realizations for inverse modeling
- Are FEM an ultimate answer?



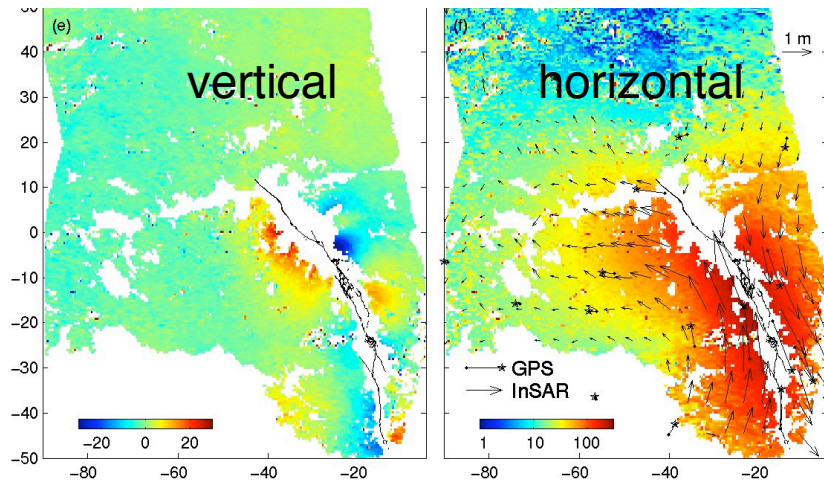


→  $\varepsilon(t) = F(\sigma, E, t, D, \dots)$  →

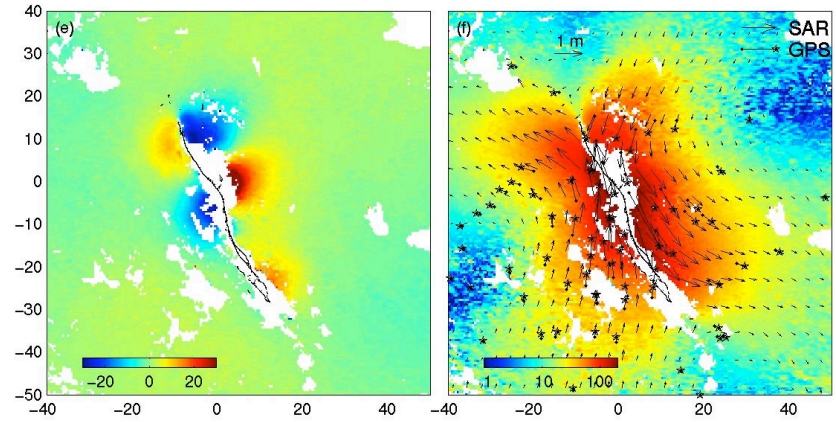




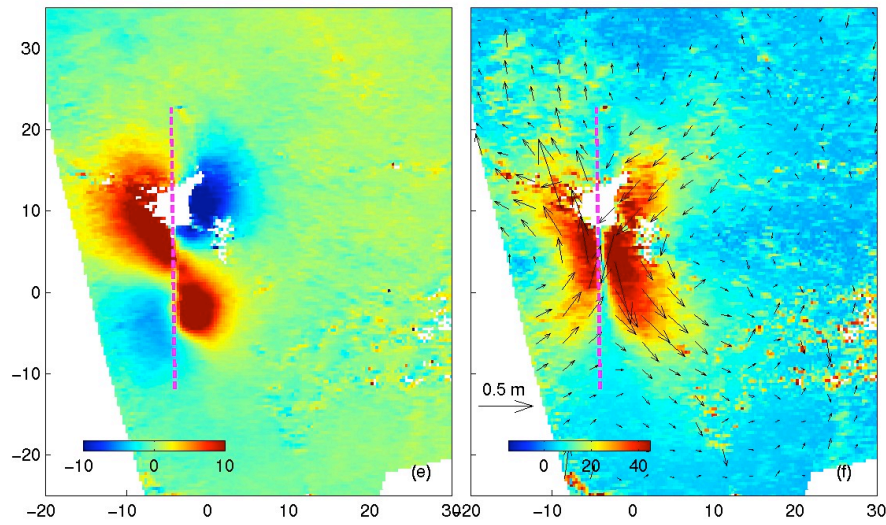
M7.3 Landers, 1992



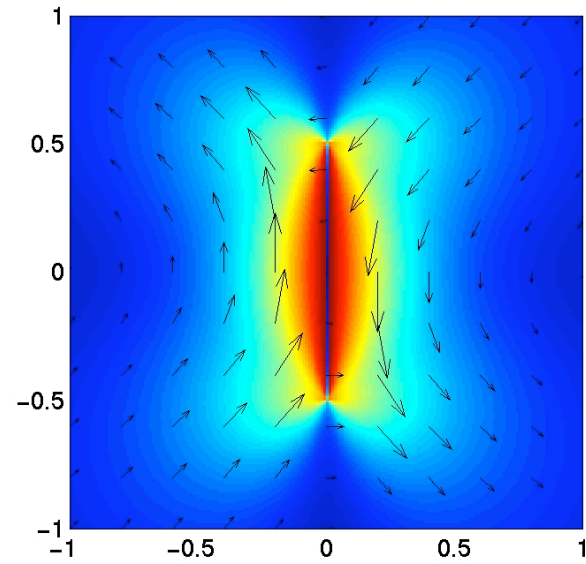
M7.1 Hector Mine, 1999

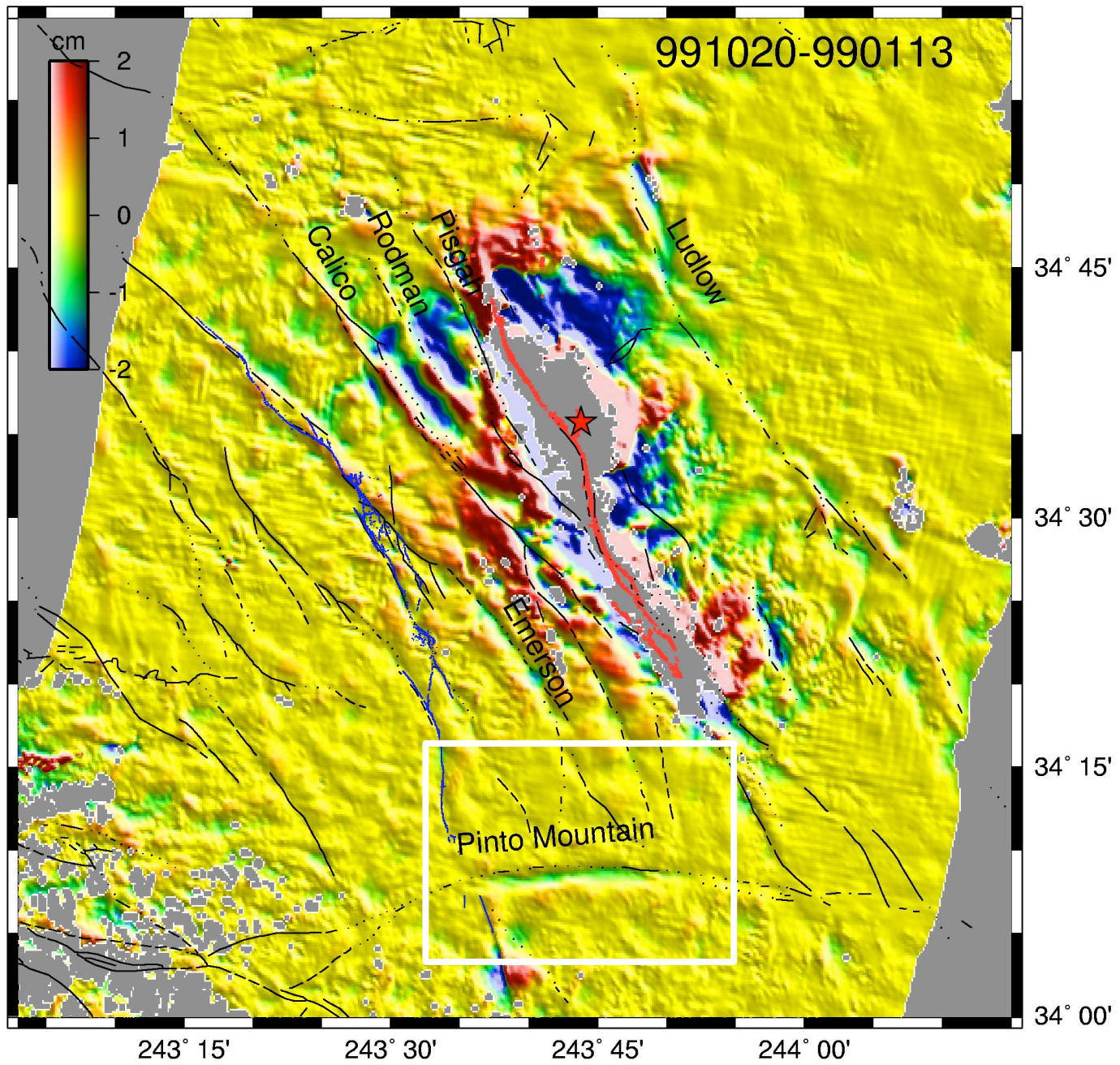


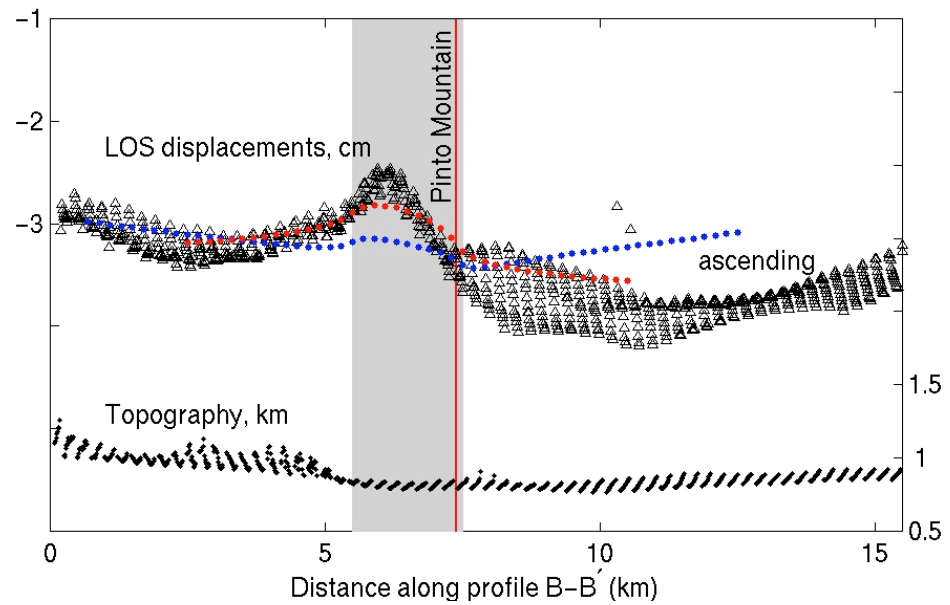
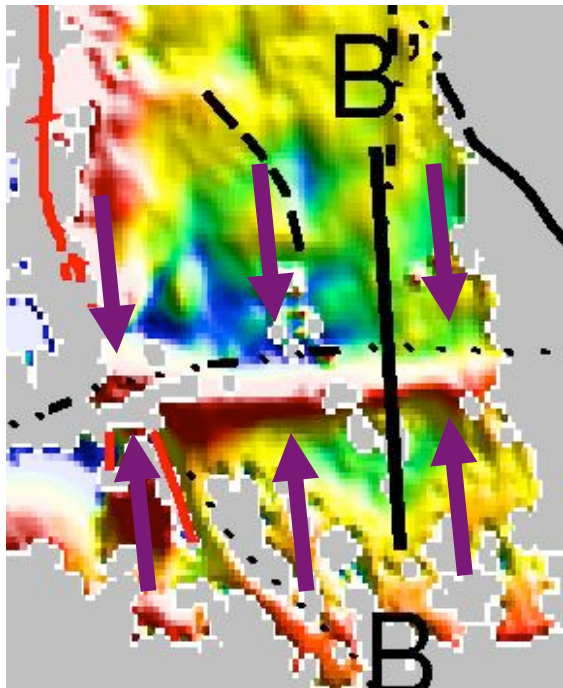
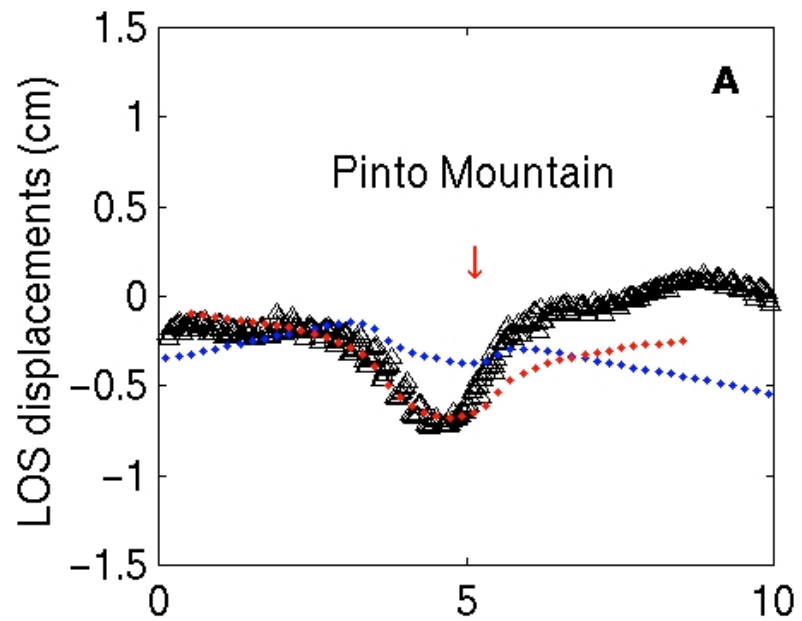
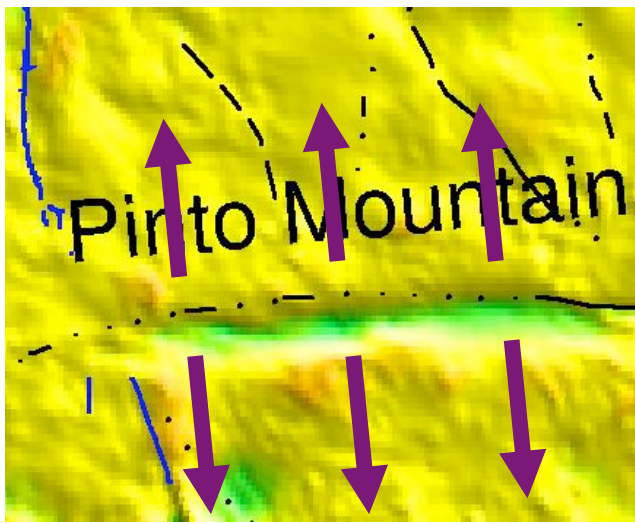
M6.6 Bam (Iran), 2003



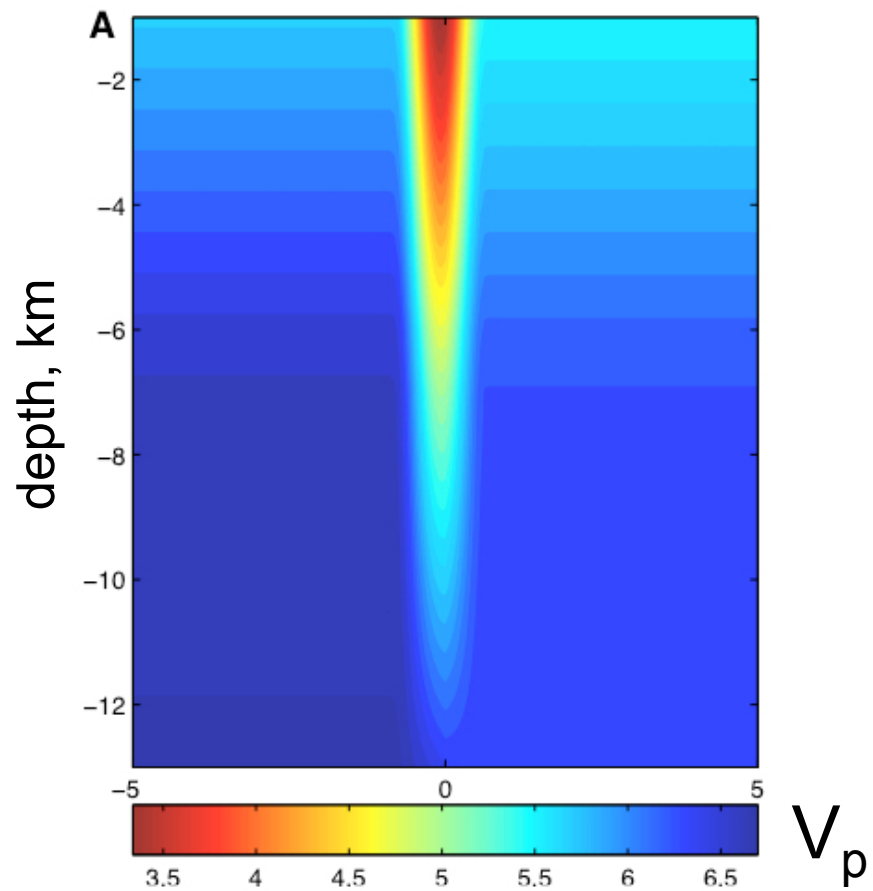
Dislocation



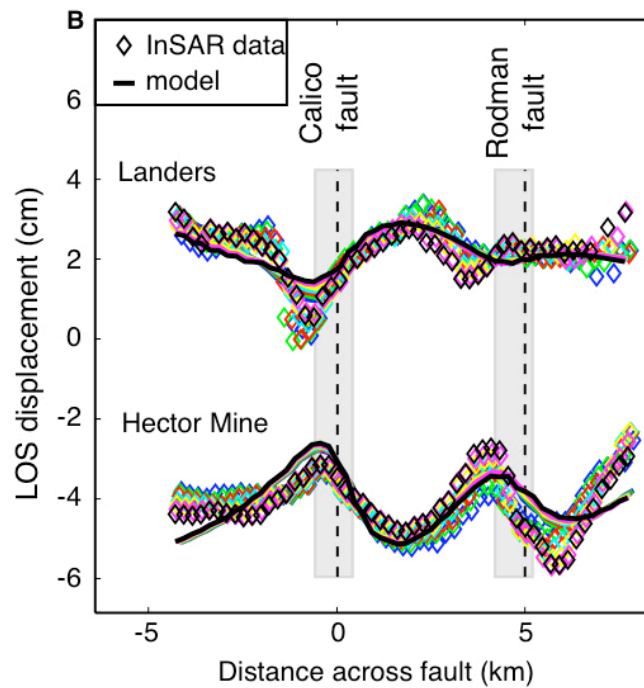




# Calico fault



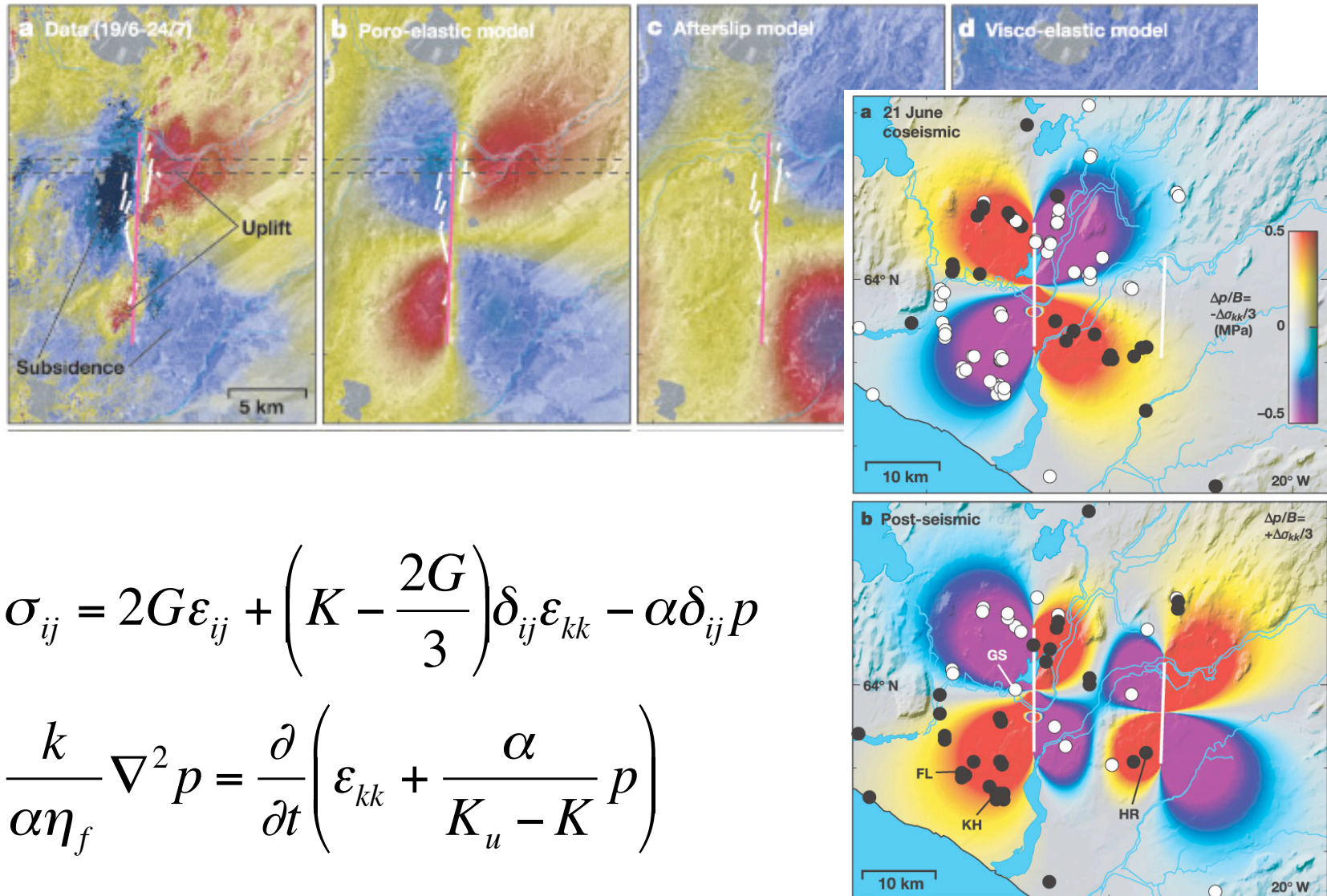
seismic tomography

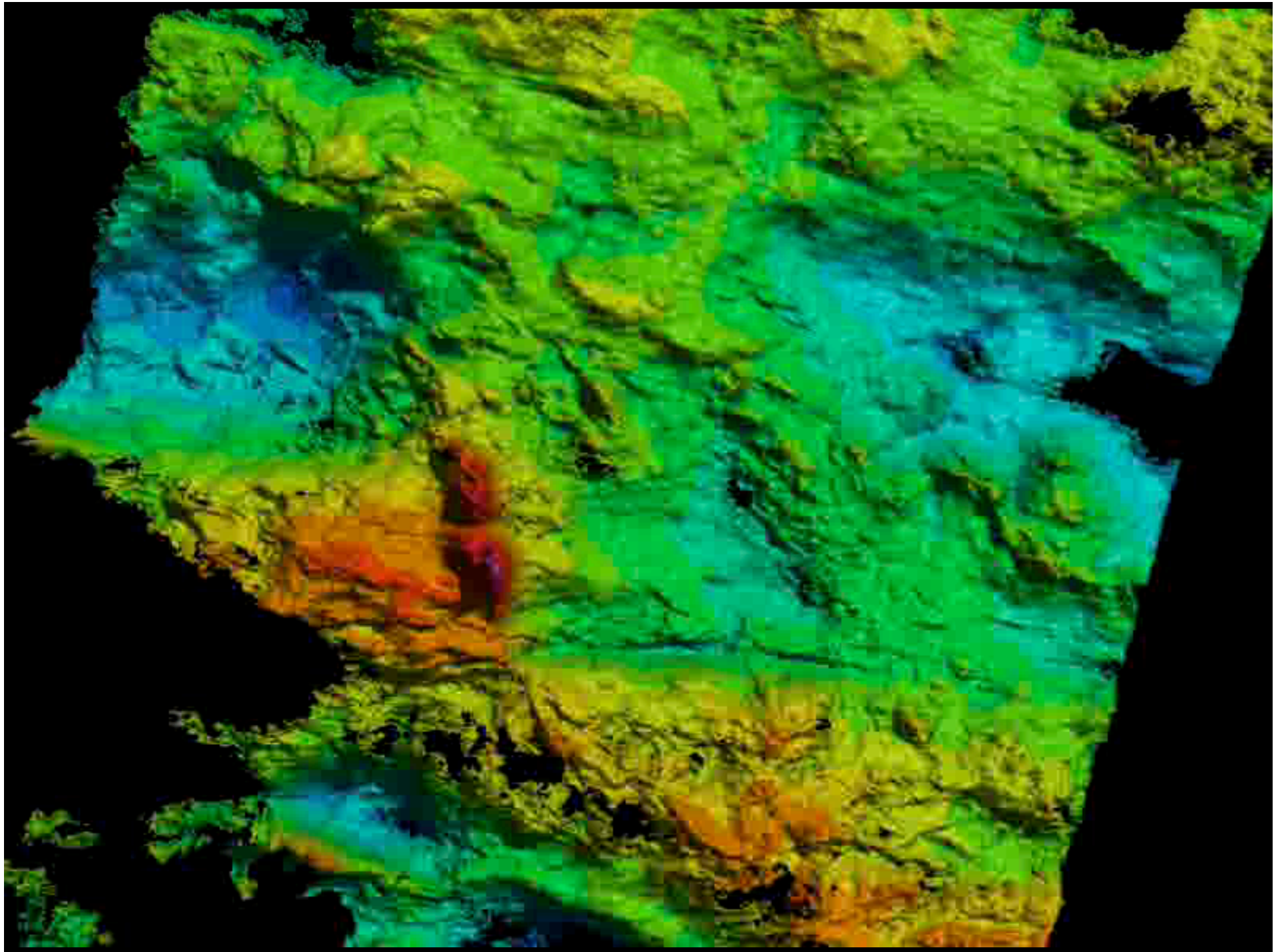


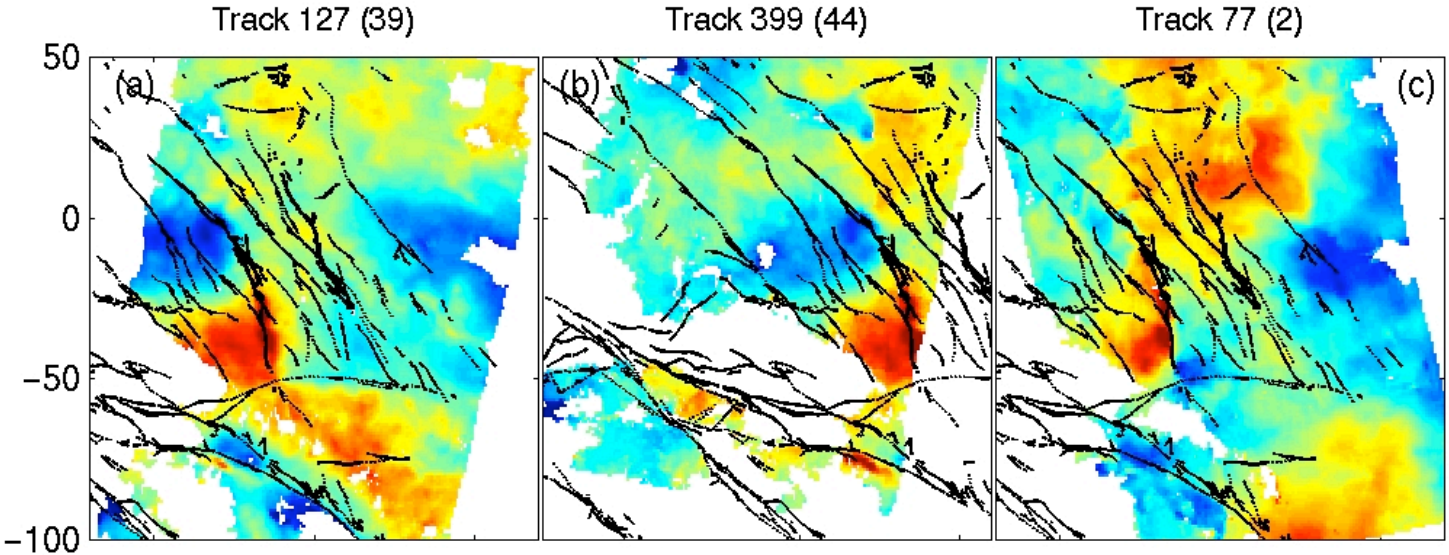


# Time-dependent deformation following earthquakes: Common suspects

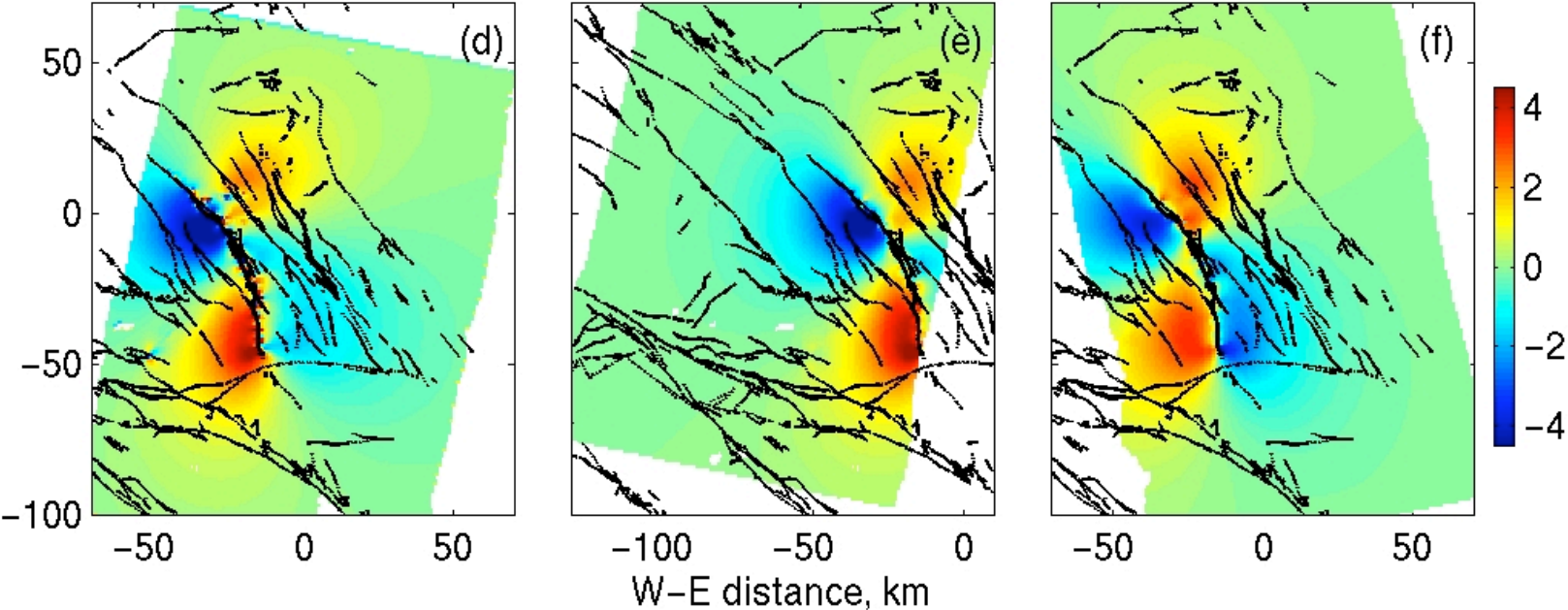
- Localized slip on or below the seismic rupture (“afterslip”)
- Visco-elastic relaxation (lower crust/upper mantle; various stress-strain relationships)
- Poro-elastic rebound (incapable of large horizontal displacements; mostly vertical deformation)
- ...or a combination of the above







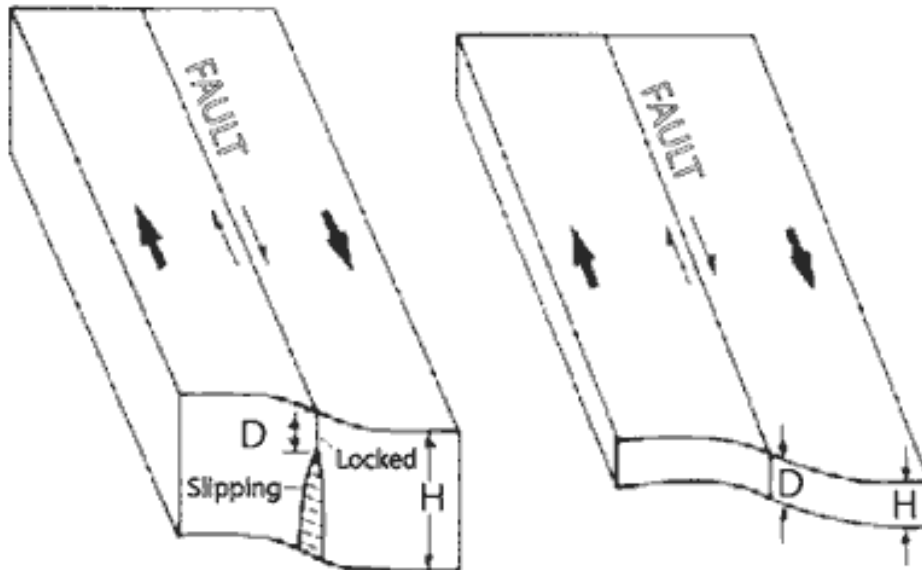
Post-seismic deformation due to the M7.3 Landers earthquake



Thick  
Lithosphere  
Model  
 $D/H \ll 1$

VS

Thin  
Lithosphere  
Model  
 $D/H = 1$

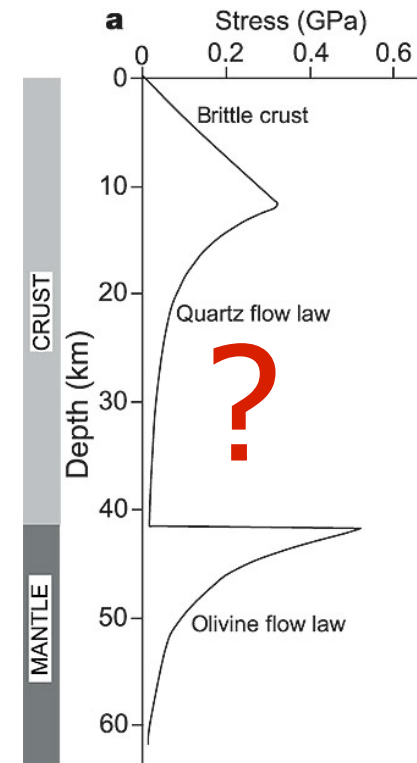


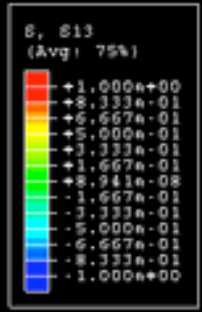
Savage and Burford, 1970  
Thatcher, 1983

Elsasser, 1969  
Savage and Prescott, 1978

“Thin viscous sheet vs Fault-block”

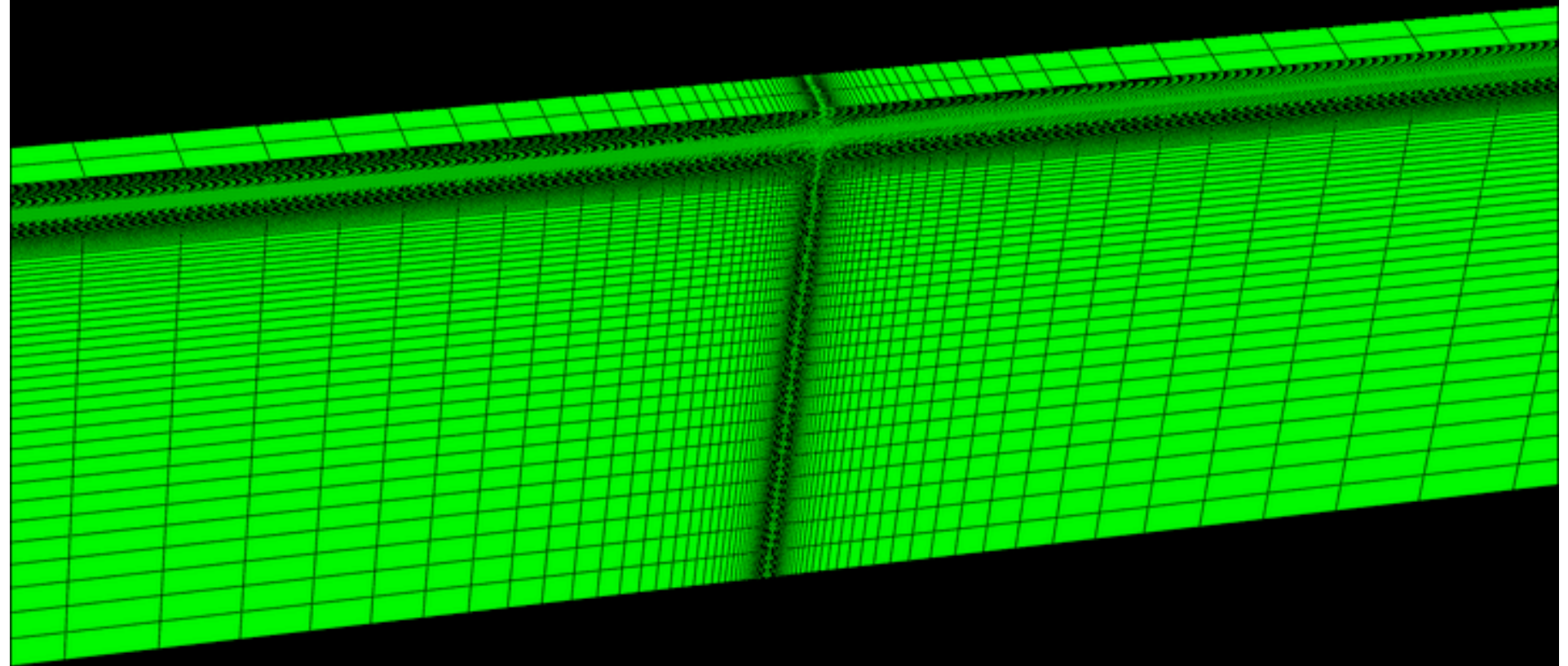
“Jelly Sandwich vs Crème Brule”





V=4 cm/yr

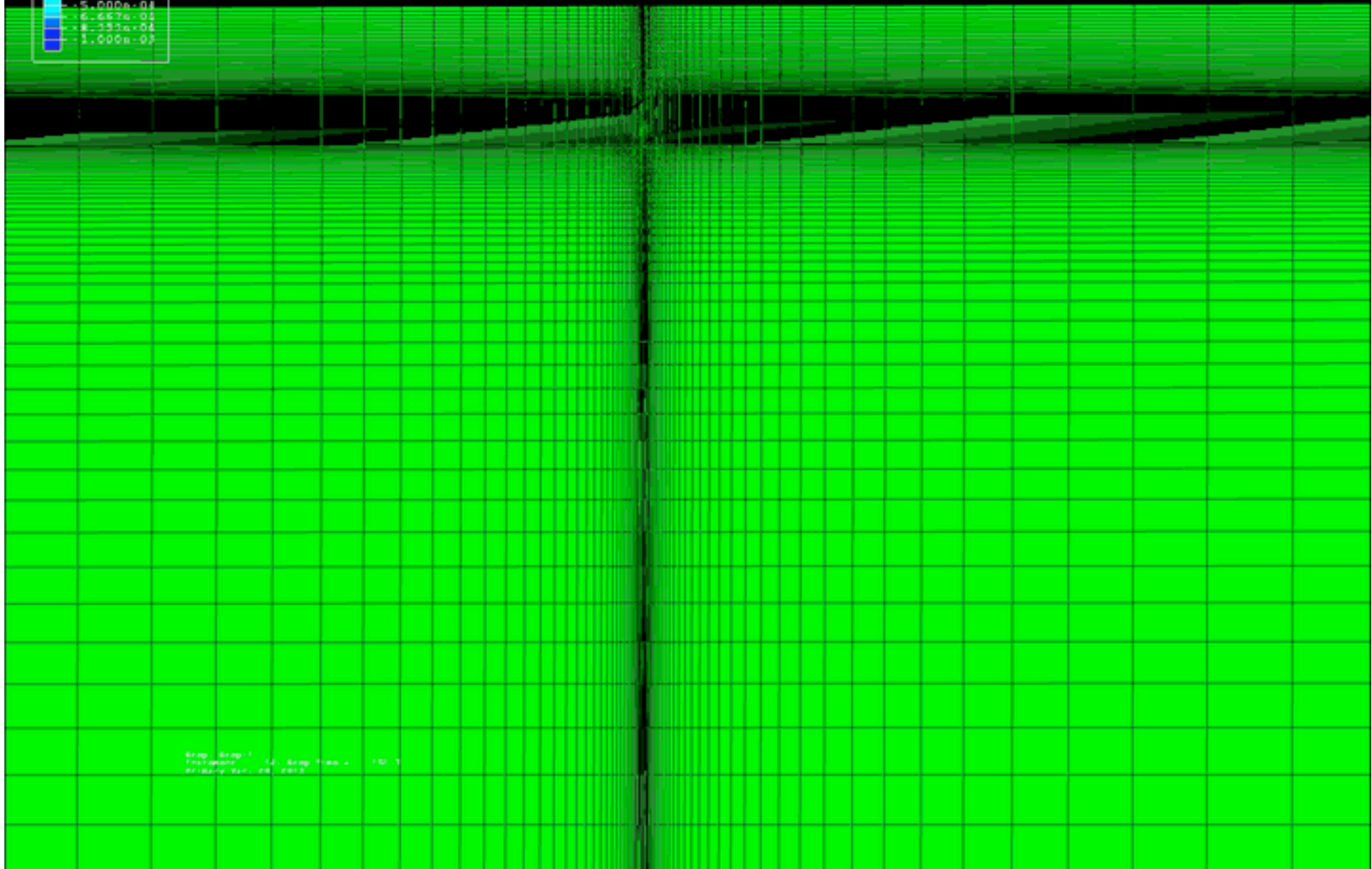
$t_r=200$  yrs  
D=8 m



Step: Step-1  
Increment: 0; Step Time = 0.000  
Primary Var:  $\epsilon, \epsilon_{13}$   
Deformed Var: U Deformation Scale Factor: +3.000e+03



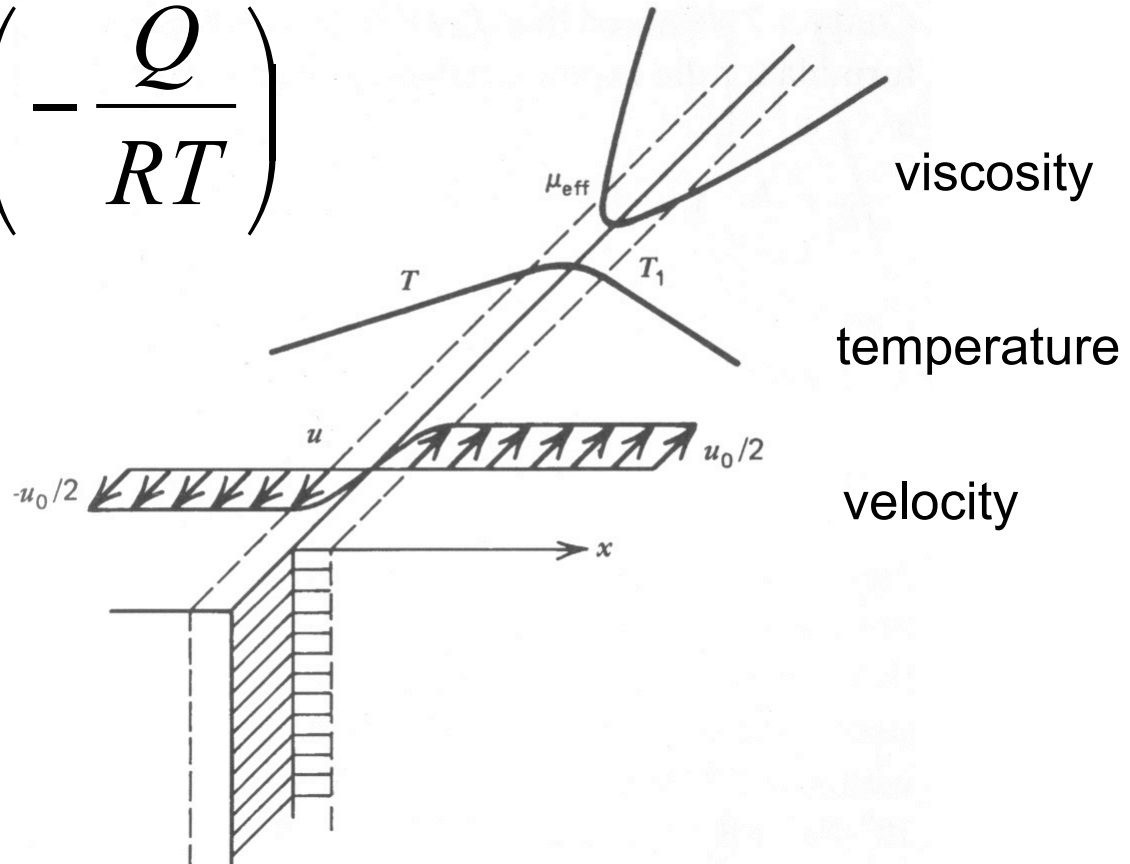
$$\dot{\epsilon} = A\sigma^{3.5}$$



## Thermo-mechanical coupling

$$\dot{\varepsilon} = C \sigma^n \exp\left(-\frac{Q}{RT}\right)$$

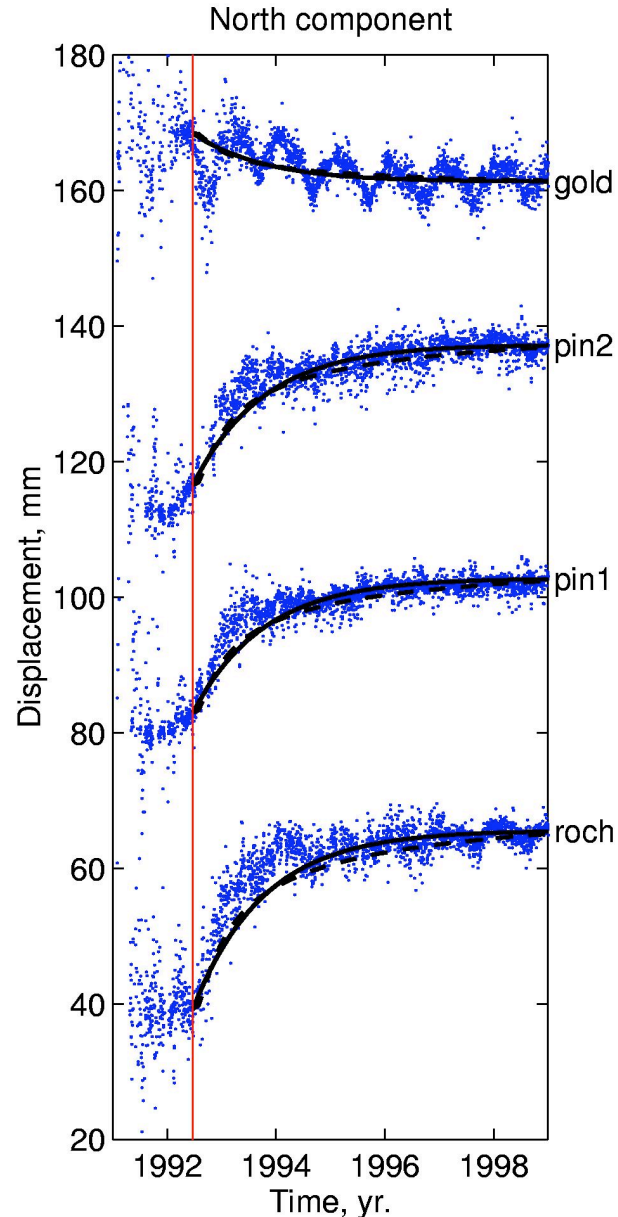
$$T \propto \dot{\varepsilon} \sigma$$



Yuen et al., 1978; Fleitout and Frodivaux, 1980;  
Turcotte and Schubert, 2002



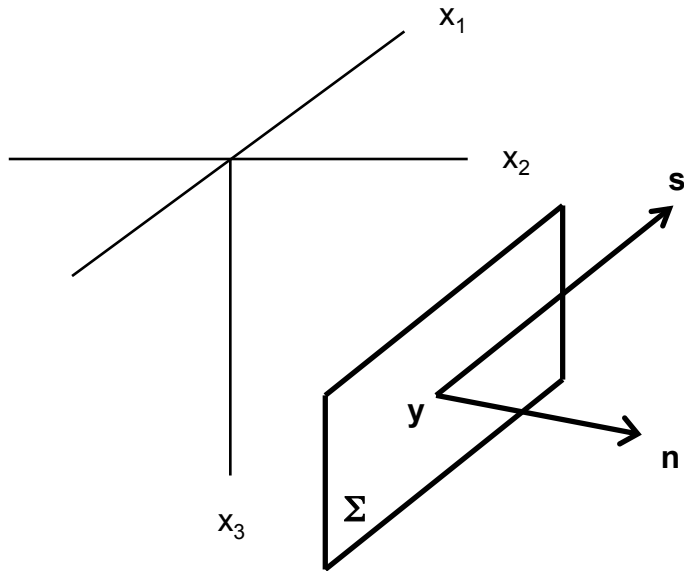
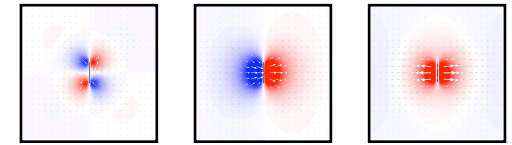
# Evidence for and implications from non-linear rheologies



Post-Landers CGPS data

- Rapid initial transient followed by a more gradual decay
- Difficult to fit assuming exponential dependence (not consistent with linear Maxwell viscoelastic behavior)
- Possible explanations:
  - Bi- (or multi-) viscous rheology
  - Power-law rheology
  - Rate-and-state friction (or some other form of non-linear localized creep)

# Equivalence between dislocations and body force couples (point-source solution)



Potency Tensor (Eigenstrain)

$$\boldsymbol{\varepsilon}^i(\mathbf{x}, \mathbf{y}) = \frac{1}{2} (\hat{\mathbf{n}} \otimes \mathbf{s} + \mathbf{s} \otimes \hat{\mathbf{n}}) \delta(\mathbf{x} - \mathbf{y})$$

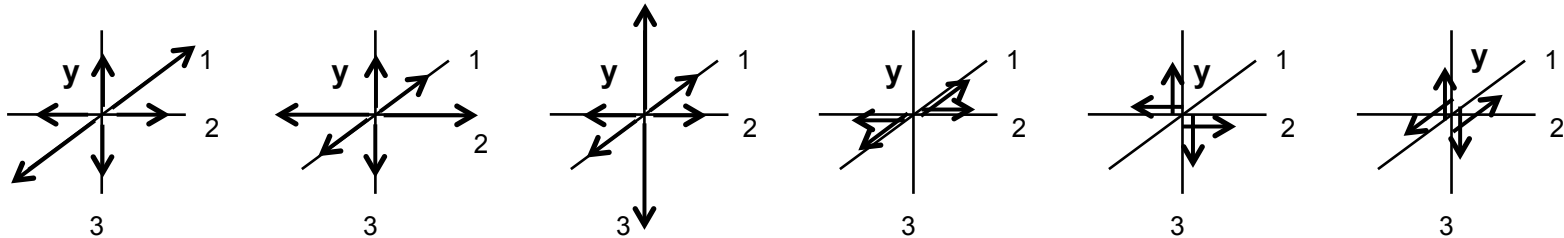
Moment Density Tensor

$$\mathbf{m}(\mathbf{x}, \mathbf{y}) = \mathbf{C} : \boldsymbol{\varepsilon}^i = \mathbf{C} : \mathbf{s} \otimes \hat{\mathbf{n}} \delta(\mathbf{x} - \mathbf{y})$$

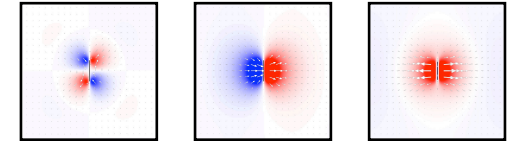
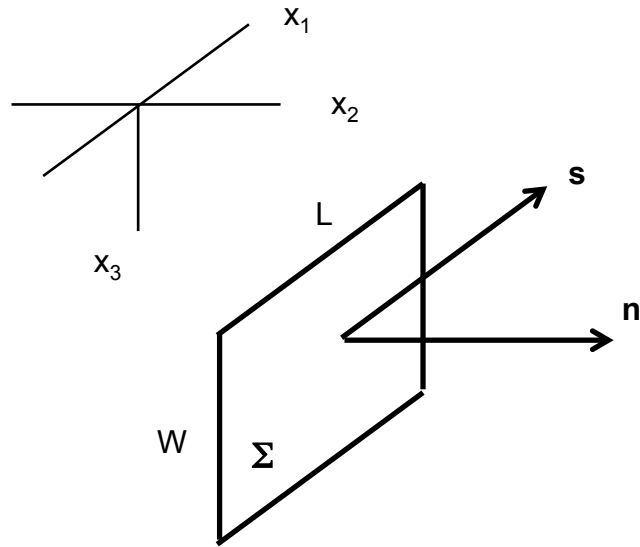
Equivalent Body Forces

$$\mathbf{f}(\mathbf{x}, \mathbf{y}) = -\nabla \cdot \mathbf{m} = -\nabla \cdot (\mathbf{C} : \boldsymbol{\varepsilon}^i)$$

Equivalent Body Forces Are Linear Combination of 6 Double Couples



# Finite Fault Source



Potency Tensor (Eigenstrain)

$$\boldsymbol{\varepsilon}^i(\mathbf{x}) = \int_{\Sigma} \boldsymbol{\varepsilon}^i(\mathbf{x}, \mathbf{y}) \, d\mathbf{y}$$

Moment Density Tensor

$$\mathbf{m}(\mathbf{x}) = \int_{\Sigma} \mathbf{C} : \boldsymbol{\varepsilon}^i(\mathbf{x}, \mathbf{y}) \, d\mathbf{y}$$

Example for Uniform Rectangular Fault

$$\boldsymbol{\varepsilon}^i(\mathbf{x}) = \frac{1}{2} (\hat{\mathbf{e}}_1 \otimes \hat{\mathbf{e}}_2 + \hat{\mathbf{e}}_2 \otimes \hat{\mathbf{e}}_1) \Pi\left(\frac{x_1 - y_1}{L}\right) \delta(x_2 - y_2) \Pi\left(\frac{x_3 - y_3}{W}\right)$$

$$\mathbf{f}(\mathbf{x}) = -2\mu \nabla \cdot \boldsymbol{\varepsilon}^i = -\mu \begin{pmatrix} \Pi\left(\frac{x_1 - y_1}{L}\right) \frac{\partial}{\partial x_2} \delta(x_2 - y_2) \Pi\left(\frac{x_3 - y_3}{W}\right) \\ \frac{\partial}{\partial x_1} \Pi\left(\frac{x_1 - y_1}{L}\right) \delta(x_2 - y_2) \Pi\left(\frac{x_3 - y_3}{W}\right) \\ 0 \end{pmatrix}$$

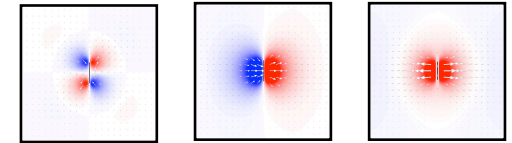
**Eigenstrain** characterizes:

- slip system (tensor part)
- location
- dimension

analytic expression for **equivalent body forces** allows:

- numerical sampling & processing
- analytic Fourier transform
- continuum representation of a discontinuous field

# Greens' Function in Fourier Domain



Navier's Equation in Space Domain

$$\nabla \cdot (\mathbf{C} : \nabla \otimes \mathbf{u}) + \mathbf{f} = 0$$

or, for isotropic elasticity

$$(\lambda + \mu) \nabla \nabla \cdot \mathbf{u} + \mu \nabla^2 \mathbf{u} + \mathbf{f} = 0$$

solution is

$$\mathbf{u} = \int_{\Omega} \mathbf{G}(\mathbf{x}, \mathbf{x}_0) \cdot \mathbf{f}(\mathbf{x}_0) dV$$

Navier's Equation in Fourier Domain

$$\mathbf{k} \cdot (\mathbf{C} : \mathbf{k} \otimes \mathbf{u}) = \mathbf{f}/4\pi$$

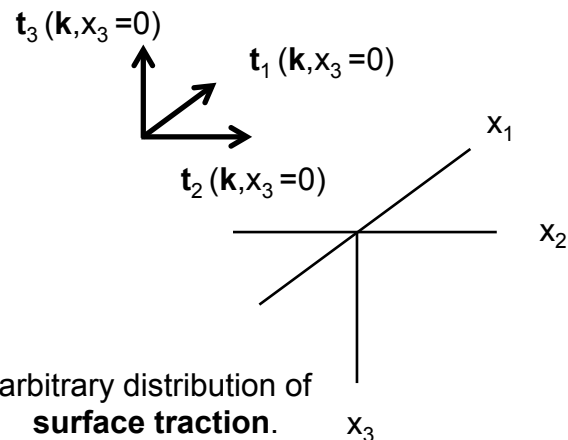
or

$$(\lambda + \mu) \mathbf{k} \otimes \mathbf{k} \cdot \mathbf{u} + \mu \mathbf{k} \cdot \mathbf{k} \otimes \mathbf{u} = \mathbf{f}/4\pi$$

or simply, with full space elastic Greens' function

$$\mathbf{G}^{-1}(\mathbf{k}; \lambda, \mu) \cdot \mathbf{u}(\mathbf{k}) = \mathbf{f}(\mathbf{k})$$

Boussinesq's & Cerruti's Problems: Elastic Deformation for Surface Traction

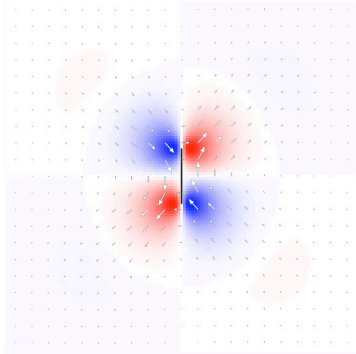


$$\mathbf{u} = \begin{pmatrix} -2B_1\beta^2 + \alpha\omega_1(B_1\omega_1 + B_2\omega_2)(1 + \beta x_3) + \alpha i\omega_1\beta B_3(1 - \alpha^{-1} + \beta x_3) \\ -2B_2\beta^2 + \alpha\omega_2(B_1\omega_1 + B_2\omega_2)(1 + \beta x_3) + \alpha i\omega_2\beta B_3(1 - \alpha^{-1} + \beta x_3) \\ \alpha \beta^2 (i(B_1\omega_1 + B_2\omega_2) x_3 - B_3(\alpha^{-1} + \beta x_3)) \end{pmatrix} e^{-\beta x_3}$$

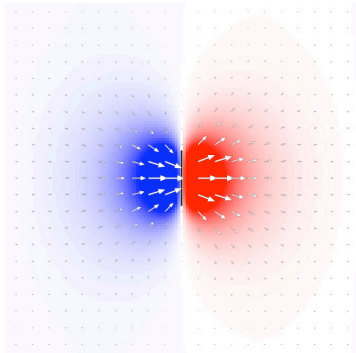
Use **Boussinesq** and **Cerruti's** solution to **remove stress** at the surface

# Benchmark

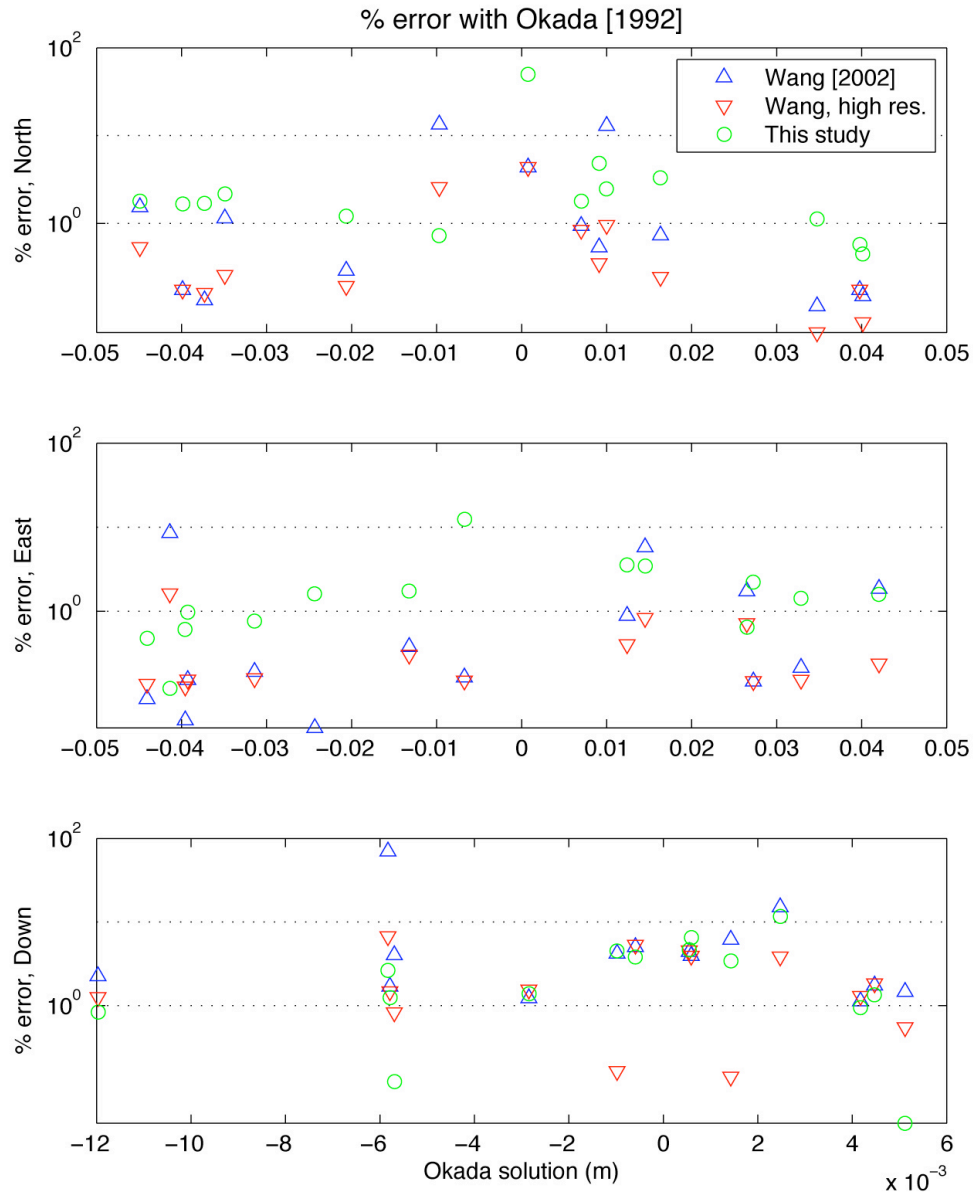
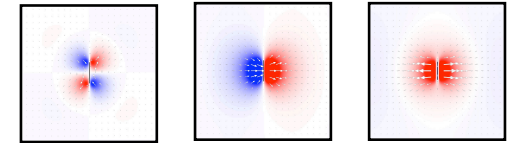
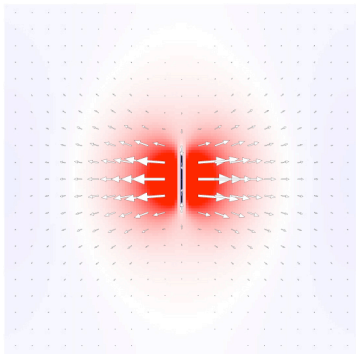
strike slip



dip slip



tensile dislocation

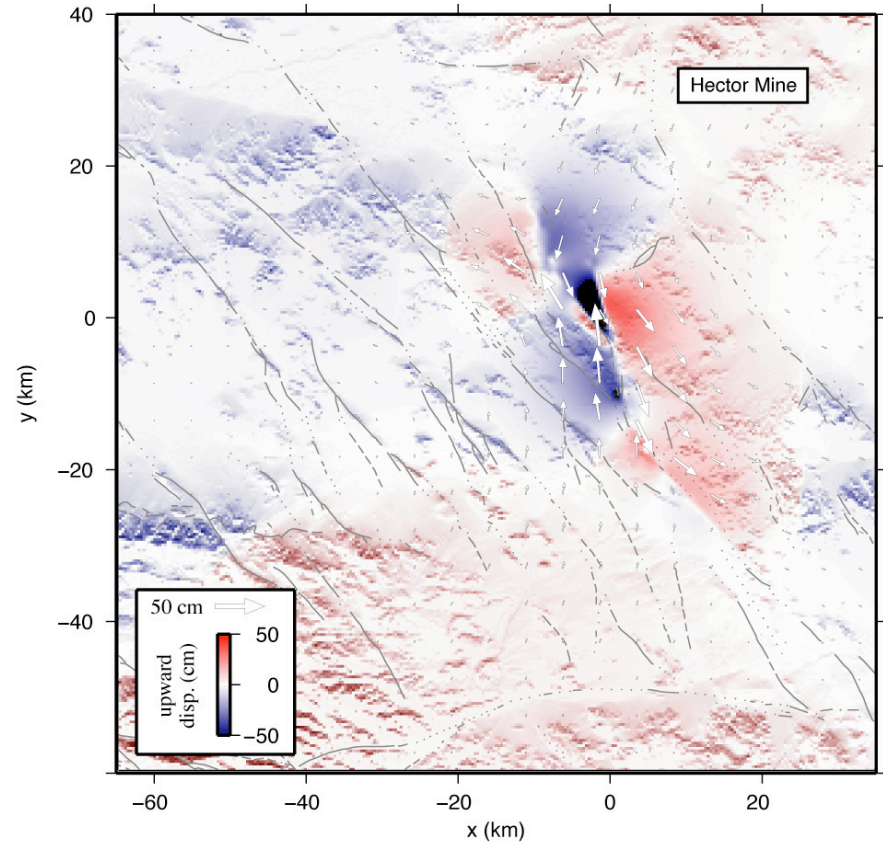
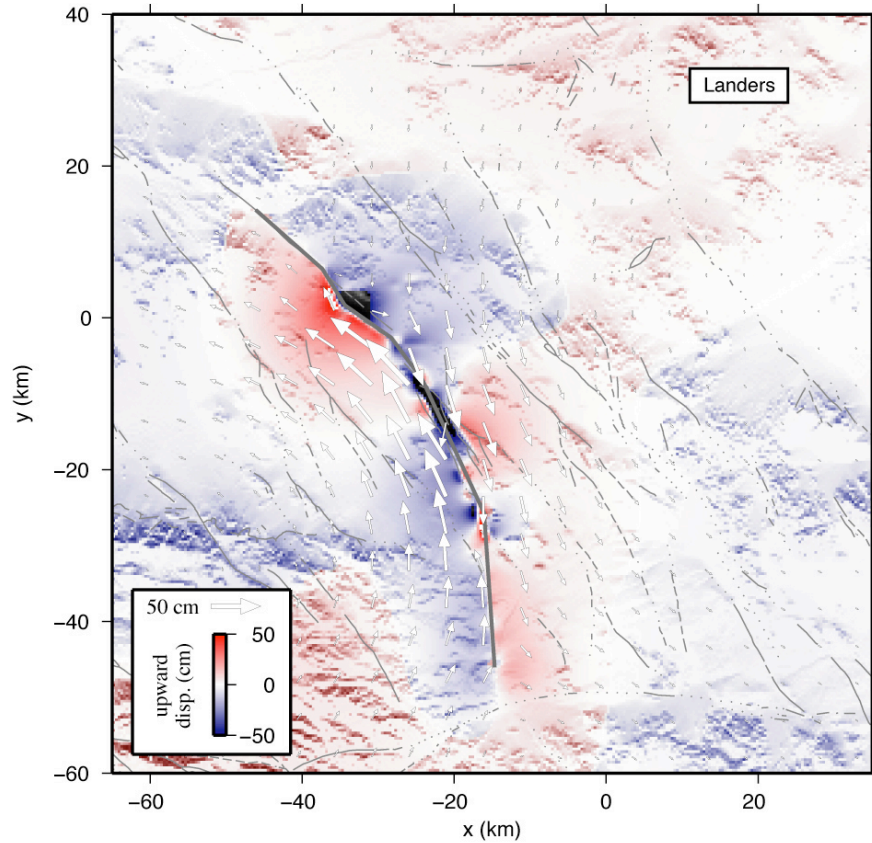
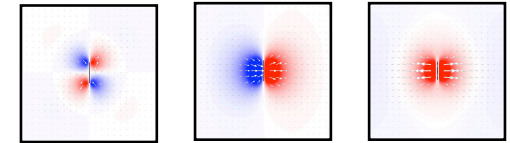


case of strike-slip fault, comparison with Okada [1992] and Wang [2002]:

- less than **5% error wrt Okada**
- comparable with Wang
- larger error in the near field (due to discontinuity approximation)

Numerical code implements **strike-slip** and **dip-slip** fault and opening (or closing) **cracks** of **arbitrary orientation**.

# Examples



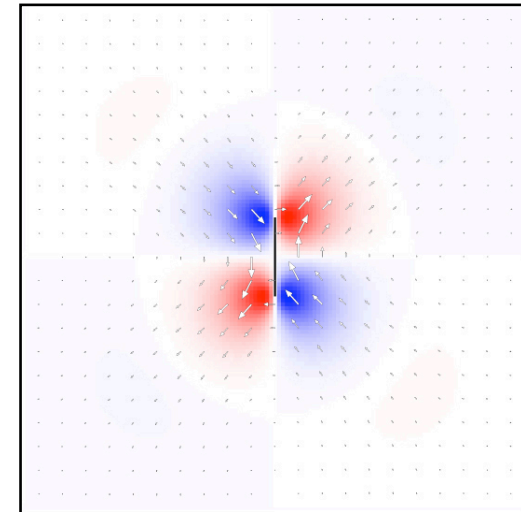
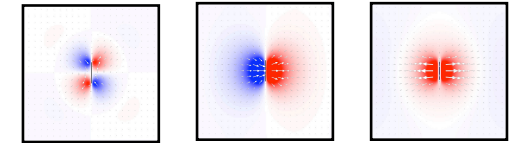
The Fourier domain method is an attractive alternative to FEM in a number of applications:

- 3-D static deformation
- nonlinear 3-D viscoelasticity
- rate-and-state fault creep
- poroelasticity
- ...

A number of mathematical issues that arise from this formulation (optimal choice of an initial “homogenized” model, convergence, existence, stability analysis, errors, etc)

# User Interface Example

```
./static <<EOF
# grid dimension (sx1,sx2,sx3)
128 128 128
# sampling (dx1,dx2,dx3), beta
0.05 0.05 0.05 0.3
# origin position (x0,y0)
0 0
# observation plane depth
0
# output directory
output
# elastic moduli (lambda,mu)
1 1
# observation points
1
1 GPS1 0.5 0.1 0
# shear dislocations
1
# index slip    x1 x2 x3 length width strike dip rake
      1      1 -0.5  0  0      1      1      0  90  0
# tensile cracks
0
EOF
```



simple interface produces output in:

- prescribed points (GPS)
- map view txt file (x1,x2,u1,u2,u3)
- Generic Mapping Tools (GMT)

code is implemented in Fortran90  
and uses DFT fourt

above example runs in 10s on a low-end  
laptop