

STRUCTURE OF THE LEEDS DYNAMO CODE

Ashley P. Willis

a.p.willis@shef.ac.uk

Governing equations

$$\begin{aligned}(Ro \partial_t - E \nabla^2) \mathbf{u} &= \mathbf{N}_u - \nabla \hat{p}, \\ (\partial_t - \nabla^2) \mathbf{B} &= \mathbf{N}_B, \\ (\partial_t - q \nabla^2) C &= N_C.\end{aligned}$$

Toroidal-poloidal decomposition

$$\begin{aligned}\mathbf{u} &= \nabla \times (T\mathbf{r}) + \nabla \times \nabla \times (P\mathbf{r}), \\ \mathbf{B} &= \nabla \times (\mathcal{T}\mathbf{r}) + \nabla \times \nabla \times (\mathcal{P}\mathbf{r}).\end{aligned}$$

Numerical formulation

$$\begin{aligned}(\partial_t - \nabla^2) \mathcal{P} &= N_{\mathcal{P}} \\ (\partial_t - \nabla^2) \mathcal{T} &= N_{\mathcal{T}} \\ (Ro \partial_t - E \nabla^2) T &= N_T, \\ (Ro \partial_t - E \nabla^2) P &= g \\ -\nabla^2 g &= N_P \\ (\partial_t - q \nabla^2) C &= N_C.\end{aligned}$$

Model equation

$$(a \partial_t - b \nabla^2) f = N$$

- A lot of code-sharing.
- narrow f.d. stencil or higher-order approximation for a given width.
- good stability.
- Each 2nd-order + 2 BCs, except system

$$\begin{cases} (Ro \partial_t - E \nabla^2) P = g \\ -\nabla^2 g = N_P \end{cases} \quad (r_i < r < r_o) \quad P = \partial_r P = 0 \quad (r = r_i, r_o),$$

which has 4 BCs on P , none on g .

Green's function or influence-matrix method

$$\begin{cases} (Ro \partial_t - E \nabla^2)P = g \\ -\nabla^2 g = N_P \end{cases} \quad (r_i < r < r_o) \quad P = \partial_r P = 0 \quad (r = r_i, r_o).$$

Write as system

$$\begin{cases} \mathbf{X} P = g \\ \mathbf{Y} g = N \end{cases} \quad P = 0, 0 \quad \partial_r P = 0, 0 \quad (r = r_i, r_o).$$

Let $P = \bar{P} + a P_i + b P_o$, where

$$\begin{cases} \mathbf{X} P_i = g_i & \partial_r P_i = 0, 0 \\ \mathbf{Y} g_i = 0 & g_i = 1, 0 \end{cases}$$

$$\begin{cases} \mathbf{X} P_o = g_o & \partial_r P_o = 0, 0 \\ \mathbf{Y} g_o = 0 & g_o = 0, 1 \end{cases}$$

$$\begin{cases} \mathbf{X} \bar{P} = \bar{g} & \partial_r \bar{P} = 0, 0 \\ \mathbf{Y} \bar{g} = N & \bar{g} = 0, 0 \end{cases}$$

- P_i and P_o both precomputable.
- Only inverted for \bar{P} each timestep.
- P satisfies two BCs automatically: $\partial_r P = 0, 0$.
- Scalars a and b set by remaining two BCs: $P = 0, 0$.
(Requires inversion of a 2×2 matrix.)
- BCs on another variable or coupling BCs satisfied to numerical precision.
- Computational overhead is nominal.

Modules

Somewhere to [gather together related data and functions](#) that work on that data. Don't worry about syntax here!

```
!*****
module rotation
!*****
  use parameters, timestep      ! use data from other modules
  implicit none
  save

                                ! data other mods can inherit
  double precision :: rot_omega
  double precision :: rot_torque

                                ! private data
  double precision, private :: rot_inertia

contains
!-----
!  initialise
!-----

  subroutine rot_precompute()
    rot_omega = 0d0
    rot_inertia = (8d0*d_PI/15d0) * d_ICradius**5
  end subroutine rot_precompute

!-----
!  timestep
!-----

  subroutine rot_predictor()
    rot_omega = rot_omega &
      + (tim_dt/rot_inertia) * rot_torque
  end subroutine rot_predictor

!*****
end module rotation
!*****
```

Modules:

codensity, (IC) rotation, velocity, magnetic	these variables stored here + predictor-corrector functions
timestep	functions for setting up matrices for model eqn $(a\partial_t - b\nabla^2) f = N.$
nonlinear	evaluate nonlinear (coupling) terms $N.$
mpi	mpi_rnk, mpi_size
mesh	selection of radial points r_n
legendre	weights for transforms.
variables	definition of data types (coll), (spec), (phys)
transform	between spherical harmonic coeffs and data at points in physical space.

The 'main' program:

parameters	physical and numerical parameters
main	the main time-stepping loop
io	input / output

A 'derived' data type:

Define type:

```
type phys                                !theta,phi,r
  double precision :: Re(10, 0:11, 5)
end type phys
```

Declare a variable: type (phys) :: p

Set value of an element: p%Re(j,i,n) = 1d0

- No need to specify the dimension of every variable with the same size
- Very handy when passing between functions.
- Fortran column dominant – keep 'close' data in first index, split over last index...

Legendre transform strongly linked to data types

$$A(\theta, \phi) = \sum_{m=0}^{M-1} \sum_{l=m}^{L-1} A_{lm} \hat{Y}_l^m(\theta, \phi)$$

$$A(r_n, \theta, \phi) = \sum_m e^{im\phi} \sum_l A_{lmn} P_l^{|m|}(\cos \theta)$$

$$A(r_n, \theta_j, \phi) = \sum_m e^{im\phi} A_{jmn}$$

$$A(r_n, \theta_j, \phi_i) = A_{jin}$$

(coll) $(A_n)_{lm}$ Data at collocation point for each harmonic.
All n on same CPU (split accross modes lm)

→ transpose →

(spec) $(A_{lm})_n$ All spectral coeffs on same CPU (split accross radial pts)

sum over l at each θ_j → $(A_{jm})_n$

sum over m at ϕ_i (FFT) →

(phys) $(A_{ij})_n$ data evaluated on points in physical space.

Strengths

- **Short** – 3000 lines excluding IO. Readable – modular, function-based.
- Data easily manipulated – function-based, netCDF.
- **Fast** in serial.
- Timestep control (improved by Chris Jones).
- **Parallelised**, linear scaling with number of CPUs.

Weakness

- Number of **CPUs currently limited** to number of radial points.
(See index n on (spec))

Possible way out:

$(A_n)_{lm}; (\text{transpose}) \rightarrow (A_l)_{mn}; (\text{sum over } l) \rightarrow (A_j)_{mn};$
 $(\text{transpose}) \rightarrow (A_m)_{jn}; (\text{FFT}) \rightarrow (A_i)_{jn} .$

The 'derived' types:

```
type spec
  double precision      :: Re(0:i_H1, i_pN)
  double precision      :: Im(0:i_H1, i_pN)
end type spec

type coll
  double precision      :: Re(i_N, 0:i_pH1)
  double precision      :: Im(i_N, 0:i_pH1)
end type coll

type phys
  double precision      :: Re(i_Th, 0:i_Ph-1, i_pN)
end type phys
```

Currently split over radial points, $n \in [1, N]$,
and harmonics lm mapped to a single index $nh \in [0 : H - 1]$.

```
nh = -1
do m = 0, M-1
  do l = m, L-1
    nh = nh + 1
```